# A Redetermination of the Crystal and Molecular Structure of Nickel(II) Ethylmethylglyoxime* 

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#### Abstract

The crystal and molecular structure of bis-(2,3-pentanedionedioximate- $N, N^{\prime}$ ) nickel(II) [or nickel(II) ethylmethylglyoxime] has been determined from three-dimensional X-ray diffraction data. The complex crystallizes in the monoclinic system in space group $P 2_{1} / c$ with two molecules per unit cell. Unit-cell parameters are: $a=4.7471 \pm 0.0005, b=11.7409 \pm 0.0030, c=11.9895 \pm 0.0020 \AA ; \beta=91.611 \pm 0.016^{\circ}$. The structural analysis of the complex, based on all reflections out to $50^{\circ}$ in $2 \theta$ with Mo $K \alpha$ radiation (final $R$ value for 802 observed reflections of $4.8 \%$ ) reveals an intramolecular hydrogen bond with an $\mathrm{O}-\mathrm{-O}$ distance of $2.454 \AA$. Bond distances in the five-membered chelate ring indicate significant delocalization of the two conjugated carbon-nitrogen double bonds.


## Introduction

The vic-dioxime complexes of metal(II) ions have received considerable attention, not only because they are useful in the analysis of certain transition metals, but also because the solid complexes exhibit certain structural features which are of general interest. These features include short intramolecular hydrogen bonds and packing configurations which give rise to unusual optical properties.

The detailed study of such features and the correct explanation of these unusual optical properties are dependent on precise structural information. The structure of nickel(II) ethylmethylglyoxime [ $\mathrm{Ni}(\mathrm{EMG})_{2}$ ] was originally solved (Frasson \& Panattoni, 1960) with two-dimensional film data. Infrared structural studies of the hydrogen bonding in several metal(II)-vic-dioximes by Caton \& Banks (1967), together with a theoretical treatment of hydrogen bonding by Lippincott \& Schroeder (1955), suggest that the intramolecular hydrogen bond in $\mathrm{Ni}(\mathrm{EMG})_{2}$ is considerably longer than the reported value of $2.33 \AA$. We have therefore undertaken a redetermination of the crystal and molecular structure of $\mathrm{Ni}(\mathrm{EMG})_{2}$.

## Experimental

## Crystal data:

$\mathrm{C}_{10} \mathrm{H}_{18} \mathrm{~N}_{4} \mathrm{O}_{4} \mathrm{Ni}, \quad M=316.99$ g.mole ${ }^{-1}$, Monoclinic $P 2_{1} / c, \quad a=4.7471 \pm 0.0005, \quad b=11.7409 \pm 0.0030, \quad c=$ $11.9895 \pm 0.0020 \AA, \beta=91.611 \pm 0.016^{\circ}, V=668.0 \AA^{3}$, $D_{c}=1.58 \mathrm{~g} . \mathrm{cm}^{-3}, Z=2, F(000)=332, T \simeq 24^{\circ}$, Мо $K \alpha$ $(\lambda=0.7107 \AA), \mu=14.7 \mathrm{~cm}^{-1}$.

Two different modifications of $\mathrm{Ni}(\mathrm{EMG})_{2}$ have been reported (Anex \& Krist, 1967). One of these crystallizes in a monoclinic space group and has been desig-

[^0]nated the $\alpha$ form. The other crystallizes in an orthorhombic space group and has been designated the $\beta$ form. A recent study by Egneus (1969) has shown that the $\alpha$ form is the more stable of the two and that crystallization from chloroform produces this modification alone. Crystals of $\mathrm{Ni}(\mathrm{EMG})_{2}$ were prepared by precipitation of the complex from aqueous solution followed by recrystallization from a chloroform solution. This procedure was employed in order to insure that the crystals obtained would be of the $\alpha$ form.

A suitable crystal was selected and mounted in a thin-walled Lindemann glass capillary. Preliminary precession photographs exhibited $2 / m$ Laue symmetry, indicating a monoclinic space group. Systematic absences of $h 0 l$ for $l$ odd and $0 k 0$ for $k$ odd indicated the space group $P 2_{1} / c\left(C_{2 h}^{5}\right)$ in agreement with the previously reported determination. The unit-cell parameters and their standard deviations were obtained by a least-squares fit to 14 independent reflection angles whose centers were determined by left-right, top-bottom beam splitting on a previously aligned HilgerWatts four circle diffractometer (Mo $K \alpha$ radiation, $\lambda=0.7107 \AA$ ). Any error in the instrumental zero was eliminated by centering the reflection at both $+2 \theta$ and $-2 \theta$.

For data collection, a crystal having approximate dimensions $0.2 \times 0.2 \times 0.5 \mathrm{~mm}$ was mounted in a glass capillary with a along the spindle axis. Data were collected at room temperature using a fully automated Hilger-Watts four circle diffractometer equipped with scintillation counter and interfaced with an SDS-910 computer in a real time mode. Two octants of data were recorded using Zr -filtered Mo $K \alpha$ radiation within a $2 \theta$ sphere of $50^{\circ}(\sin \theta / \lambda=0 \cdot 5947)$. The $\theta-2 \theta$ step-scan technique, $0.01^{\circ}$ per step counting for 0.4096 sec per step, was employed with a take-off angle of $4.5^{\circ}$. Variable step symmetric scan ranges were used with the number of steps for a given reflection determined as follows: $N=\left(50+2\right.$ per $\left.{ }^{\circ} \theta\right)$. Stationary-crystal station-ary-counter background measurements were made by counting for one half the total scan time at each end of
the scan. 1608 reflections were measured in this way.
As a general check on the electronic and crystal stability, the intensities of three standard reflections were measured periodically during the data collection. Monitoring options based on these standard counts were used to maintain crystal alignment and to stop the collection of data if standard counts fell below statistically allowed fluctuations. No decrease in any of the standards was observed during data collection.
The intensity data were corrected for Lorentz- polarization effects. The absorption coefficient, $\mu$, is 14.68 $\mathrm{cm}^{-1}$, and an absorption correction (Busing \& Levy, 1957) was made using $A B C O R$;* the maximum and minimum transmission factors were 0.839 and 0.755 respectively. The estimated error in each intensity was calculated by $[\sigma(I)]^{2}=\left[C_{t}+C_{b}+\left(0.05 C_{t}\right)^{2}+\left(0.05 C_{b}\right)^{2}\right]$ where $C_{t}$ and $C_{b}$ are the total count and background count respectively. The individual values of $F_{o}^{2}$ from equivalent sets were then averaged to give 1397 independent $F_{o}^{2}$ values. The estimated standard deviation in each structure factor was calculated from the mean deviation of intensity by the method of finite differences (Williams \& Rundle, 1964). The reciprocals of

* In addition to local programs written at this Laboratory, the authors wish to acknowledge the use of $A B C O R$ (Busing \& Levy, 1957), ORTEP (Johnson, 1965), ORFFE (Busing, Martin \& Levy, 1964), LCR2 (Williams, 1964), and a local modification of ORFLS (Busing, Martin \& Levy, 1962).


Fig. 1. Interatomic distances and angles in $\mathrm{Ni}(\mathrm{EMG})_{2}$.
the structure factor variances were used as weights in the least-squares refinement.

## Solution and refinement of structure

Owing to the presence of only two molecules per unit cell, the metal atoms are restricted to the centers of symmetry. The positions of the remaining atoms except hydrogen were obtained from three-dimensional electron density difference maps and were refined by fullmatrix least-squares techniques with isotropic thermal parameters to a conventional discrepancy index of

$$
R=\Sigma| | F_{o}\left|-\left|F_{c}\right|\right| / \Sigma\left|F_{o}\right|=0 \cdot 135,
$$

and a weighted $R$ index of

$$
w R=\left[\sum w\left(\left|F_{o}\right|-F_{c} \mid\right)^{2} / \Sigma\left|F_{o}\right|^{2}\right]^{1 / 2}=0 \cdot 122 .
$$

The relativistic Dirac-Slater X-ray scattering factors for neutral atoms of Cromer \& Waber (1965) were used with those of nickel modified for the real and imaginary parts of anomalous dispersion (Cromer, 1965). All unique data were used in the refinement.

At this point, anisotropic refinement of nickel and nitrogen atoms for two cycles gave a value for $R$ of $0 \cdot 115$. The remaining atoms with the exception of the methyl carbon atoms, were next refined anisotropically, and two cycles gave values for $R$ and $w R$ of $0 \cdot 104$ and 0.094 respectively.

An electron density difference map was calculated and it revealed some rather diffuse peaks which could be interpreted as hydrogen atoms. Owing to the uncertainty in the positions of these atoms, however, their positions were calculated assuming tetrahedral methyl groups and a C-H distance of $1.08 \AA$. The methyl hydrogen atoms were placed as near as possible to the observed peaks in the difference map. The methylene hydrogen atoms were placed in the plane which bisects the $\mathrm{C}-\mathrm{C}-\mathrm{C}$ angle of the ethyl substituent, at an angle of $109.5^{\circ}$ from one another. The hydroxyl hydrogen atom was placed midway between the oxygen atoms. All hydrogen atoms were assigned isotropic thermal parameters slightly larger than those of the atoms to which they are bound.

Two cycles of refinement with hydrogen parameters

Table 1. Final refined positional and anisotropic thermal parameters $\left(\times 10^{4}\right)$ for non-hydrogen atoms
The $\beta_{i j}$ are defined by: $T=\exp \left[-\left(h^{2} \beta_{11}+k^{2} \beta_{22}+l^{2} \beta_{33}+2 h k \beta_{12}+2 k l \beta_{23}+2 h l \beta_{13}\right)\right]$
Numbers in parentheses are standard deviations in the least significant digits.

|  | $x$ | $y$ | $z$ | $\beta_{11}$ | $\beta_{22}$ | $\beta_{33}$ | $\beta_{12}$ | $\beta_{13}$ | $\beta_{23}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| Ni | 0 | 0 | 0 | 410 (5) | 74 (1) | 53 (1) | -21 (2) | 30 (1) | -3 (1) |
| $\mathrm{N}(1)$ | -270 (8) | -1524 (4) | -414 (3) | 331 (20) | 79 (4) | 67 (3) | -13(7) | 37 (7) | 4 (3) |
| $\mathrm{N}(2)$ | 2686 (9) | 25 (4) | -1102 (3) | 517 (23) | 79 (4) | 60 (3) | -19 (9) | 30 (7) | 1 (3) |
| $\mathrm{O}(1)$ | -2071 (9) | -2272 (3) | 39 (3) | 696 (25) | 72 (3) | 100 (3) | -52(7) | 76 (7) | -3(3) |
| O(2) | 4150 (8) | 958 (3) | -1378 (3) | 635 (23) | 95 (3) | 80 (3) | -51 (8) | 82 (7) | -3(2) |
| C(1) | 1344 (11) | -1858 (4) | -1204 (4) | 524 (27) | 72 (4) | 66 (4) | 14 (10) | 2 (9) | -1 (3) |
| C(2) | 3101 (11) | -935 (4) | -1614 (4) | 426 (26) | 79 (5) | 58 (4) | 25 (9) | -4 (8) | -2 (3) |
| C(3) | 1408 (13) | -3058 (5) | -1631 (5) | 765 (39) | 76 (5) | 108 (5) | 18 (12) | 56 (12) | -12 (4) |
| C(4) | 5087 (12) | -1060 (4) | -2563 (4) | 515 (29) | 103 (5) | 75 (4) | 22 (11) | 49 (9) | -4 (4) |
| C(5) | 3594 (15) | -893 (6) | -3682 (5) | 920 (47) | 159 (8) | 74 (5) | 56 (16) | 17 (12) | -6 (5) |

unvaried and anisotropic thermal parameters for all heavy atoms gave values for $R$ and $w R$ of 0.089 and 0.074 respectively. Hydrogen atom positions were then recalculated and two further cycles gave values for $R$ and $w R$ of 0.088 and 0.073 respectively.

A statistical analysis of $w \Delta^{2}$ where $\Delta^{2}=\left(\left|F_{o}\right|-\left|F_{c}\right|\right)^{2}$ as a function of scattering angle and magnitude of $F_{o}$ revealed that the relative weighting scheme used was reasonable. An attempt to refine the positional and thermal parameters of the methyl hydrogen atoms was not entirely successful, but based on the refined $\mathrm{C}-\mathrm{H}$ bond distances these hydrogen atoms were relocated with a $\mathrm{C}-\mathrm{H}$ distance of $1.01 \AA$. During the final refinement, the positional and thermal parameters of all hydrogen atoms were calculated and not varied because the agreement between observed and calculated structure amplitudes seemed insensitive to these parameters. Two final cycles of refinement gave values for $R$ and $w R$ of 0.088 and 0.073 respectively. Two final cycles of refinement using 802 observed reflections (those with $F^{2}>3 \sigma F^{2}$ ) gave values for $R$ and $w R$ of 0.048 and 0.060 respectively, with no significant shifts in the final refined parameters. A final electron density difference map revealed no peaks greater than $0.4 \mathrm{e} . \AA^{-3}$. The final standard deviation for an observation of unit weight, $\left[\sum w \Delta^{2} /(N O-N V)\right]^{1 / 2}$, where $N O$ is the number of observations (1397) and $N V$ is the number of variables (88) was $1 \cdot 26$ electrons. During the final cycle, the largest shift in any parameter was less than 0.01 times its own $\sigma$.

Table 1 gives the final'refined positional and thermal
parameters for the heavy atoms along with their standard deviations as derived from the inverse matrix of the final least-squares refinement cycle. R.m.s. components of thermal displacement along the principal axes are given for the heavy atoms in Table 2. A list of all 1397 unique recorded and calculated structure amplitudes is found in Table 3.

Table 2. R.m.s. component of thermal displacement $(\AA)$ along principal axis $R$

|  | $R=1$ | $R=2$ | $R=3$ |
| :---: | :---: | :---: | :---: |
| Ni | $0 \cdot 1848$ (12) | $0 \cdot 2126$ (17) | $0 \cdot 2398$ (13) |
| $\mathrm{N}(1)$ | $0 \cdot 1764$ (60) | $0 \cdot 2322$ (61) | 0.2378 (53) |
| N(2) | $0 \cdot 2005$ (55) | $0 \cdot 2308$ (66) | $0 \cdot 2530$ (56) |
| $\mathrm{O}(1)$ | $0 \cdot 2064$ (50) | $0 \cdot 2449$ (46) | $0 \cdot 3146$ (51) |
| O (2) | $0 \cdot 1995$ (50) | $0 \cdot 2489$ (46) | $0 \cdot 3062$ (49) |
| C(1) | $0 \cdot 2189$ (93) | $0 \cdot 2223$ (69) | $0 \cdot 2476$ (71) |
| C (2) | $0 \cdot 2035$ (64) | $0 \cdot 2109$ (72) | $0 \cdot 2453$ (69) |
| C(3) | $0 \cdot 2192$ (73) | $0 \cdot 2706$ (75) | $0 \cdot 3114$ (78) |
| C(4) | $0 \cdot 2056$ (70) | $0 \cdot 2607$ (67) | $0 \cdot 2714$ (74) |
| C(5) | $0 \cdot 2299$ (74) | $0 \cdot 3036$ (84) | $0 \cdot 3521$ (86) |

## Results and discussion

Interatomic distances and angles with standard deviations are given in Table 4 and illustrated in Fig. 1. The standard deviations in the distances and angles were calculated using the variance-covariance matrix and $O R F F E$, and include errors in the lattice constants. The present study reveals a much longer hydrogen bond than was previously reported. The latest value of $2.454 \AA$ is in much better agreement with the distance


Fig. 2. Stereo illustration of a molecule of $\mathrm{Ni}(\mathrm{EMG})_{2}$ with therral ellipsoids scaled to enclose $50 \%$ probability.


Fig. 3. Stereo view, along a, of the crystal packing of $\mathrm{Ni}(E M G)_{2}$.
expected from existing correlations of OH stretching frequencies with $\mathrm{O}--\mathrm{O}$ distance．Statistical compar－ isons of chemically equivalent bonds reveal no signi－ ficant differences and suggest that the hydrogen bond is symmetrical．The bond distances in this complex
agree well with their equivalents in other oximes，vic－ dioximes and metal（II）－vic－dioxime complexes．

A significant feature of the complex is the delocaliza－ tion of the two carbon－nitrogen double bonds through the carbon－carbon bond in the chelate ring．Similar

Table 3．Observed and calculated structure amplitudes















保保 -12
-4
6
79




为



Table 4. Interatomic distances and angles
Distances

| $\mathrm{Ni}-\mathrm{N}(1)$ | 1.861 (4) $\AA$ | $\mathrm{Ni}-\mathrm{N}(2)$ | 1.862 (4) $\AA$ |
| :---: | :---: | :---: | :---: |
| $\mathrm{N}(1)-\mathrm{O}(1)$ | $1 \cdot 350$ (5) | $\mathrm{N}(2)-\mathrm{O}(2)$ | $1 \cdot 344$ (5) |
| $\mathrm{N}(1)-\mathrm{C}(1)$ | 1.295 (6) | $\mathrm{N}(2)-\mathrm{C}(2)$ | $1 \cdot 301$ (5) |
| $\mathrm{C}(1)-\mathrm{C}(2)$ | 1.462 (7) | $\mathrm{C}(4)-\mathrm{C}(5)$ | 1.512 (8) |
| $\mathrm{C}(1)-\mathrm{C}(3)$ | 1.499 (7) | $\mathrm{C}(2)-\mathrm{C}(4)$ | 1.506 (7) |
|  | $\mathrm{O}(1)-\mathrm{O}\left(2^{\prime}\right)$ | $2 \cdot 454$ (5) |  |
| Angles |  |  |  |
| $\mathrm{Ni}-\mathrm{N}(1)-\mathrm{O}(1)$ | 123.9 (3) ${ }^{\circ}$ | $\mathrm{Ni}--\mathrm{N}(2)-\mathrm{O}(2)$ | 123.9 (3) ${ }^{\circ}$ |
| $\mathrm{Ni}-\mathrm{N}(1)-\mathrm{C}(1)$ | 116.6 (3) | $\mathrm{Ni}-\mathrm{N}(2)-\mathrm{C}(2)$ | 115.9 (3) |
| $\mathrm{N}(1)-\mathrm{O}(1)-\mathrm{O}(2)^{\prime}$ | 97.2 (3) | $\mathrm{N}(2)-\mathrm{O}(2)-\mathrm{O}(1)^{\prime}$ | 97.5 (3) |
| $\mathrm{O}(1)-\mathrm{N}(1)-\mathrm{C}(1)$ | 119.5 (4) | $\mathrm{O}(2)-\mathrm{N}(2)-\mathrm{C}(2)$ | $120 \cdot 2$ (4) |
| $\mathrm{N}(1)-\mathrm{C}(1)-\mathrm{C}(2)$ | $112 \cdot 1$ (4) | $\mathrm{N}(2)-\mathrm{C}(2)-\mathrm{C}(1)$ | $112 \cdot 8$ (4) |
| $\mathrm{N}(1)-\mathrm{C}(1)-\mathrm{C}(3)$ | $123 \cdot 5$ (5) | $\mathrm{N}(2)-\mathrm{C}(2)-\mathrm{C}(4)$ | $123 \cdot 2$ (4) |
| $\mathrm{C}(3)-\mathrm{C}(1)-\mathrm{C}(2)$ | $124 \cdot 4$ (5) | $\mathrm{C}(4)-\mathrm{C}(2)-\mathrm{C}(1)$ | 123.9 (4) |
| $\mathrm{C}(2)-\mathrm{C}(4)-\mathrm{C}(5)$ | 111.7 (5) | $\mathrm{N}(1)-\mathrm{Ni}-\mathrm{N}(2)$ | 82.5 (2) |

effects have been observed in the glyoxime complexes of $\mathrm{Ni}(\mathrm{II})$ (Calleri, Ferraris \& Viterbo, 1967a; Murmann \& Schlemper, 1967), Pd(II) (Calleri, Ferraris \& Viterbo, 1967b) and $\operatorname{Pt}(\mathrm{II})$ (Ferraris \& Viterbo, 1969) and in dimethylglyoxime (Hamilton, 1961).

With the exception of the ethyl group and $\mathrm{O}(1)$ the molecule is planar. The equation of the best mean plane through the planar portion of the molecule and the deviations of all atoms from the plane are given in Table 5. The nearest $\mathrm{O}(1)$ contact ( $3.37 \AA$ ) is with the $\mathrm{C}(2)$ atom of the molecule displaced in the a direction.

Table 5. Distances of all atoms from the best mean plane through those atoms for which t values are given

| $3.284 x-2.753 y+7.952 z=0.0^{*}$ |  |  |  |
| :---: | :---: | :---: | :---: |
|  | ${ }^{d}$ | ${ }^{\sigma}$ | $t \dagger$ |
| Ni | $0.0000 \AA$ | $0.0001 \AA$ | 0.00 |
| $\mathrm{N}(1)$ | 0.0016 | 0.0041 | $0 \cdot 40$ |
| $\mathrm{N}(2)$ | -0.0009 | 0.0037 | 0.23 |
| $\mathrm{O}(1)$ | -0.0242 | 0.0035 |  |
| $\mathrm{O}(2)$ | 0.0033 | 0.0036 | 0.92 |
| C(1) | -0.0041 | 0.0051 | 0.81 |
| C(2) | -0.0078 | 0.0050 | 1.55 |
| C(3) | 0.0071 | 0.0057 | 1.26 |
| C(4) | -0.0757 | 0.0054 |  |
| C(5) | $-1.5017$ | 0.0072 |  |
|  | $\chi^{2}=5$ | r $v=4$ |  |

* The equation given is for the molecule at the origin and is given in fractional coordinates referred to the crystallographic axes.
$\dagger t=d / \sigma ; \chi^{2}=\Sigma t^{2}$.
An indication of the directions and r.m.s. amplitudes of vibration for the non-hydrogen atoms is provided by Fig. 2. A critical examination of the thermal parameters reveals nothing unusual or unexpected in the thermal motion of these atoms. The motion of the atoms in the chelate ring is more nearly isotropic than that of the atoms bound to the ring. The motion of the terminal carbon atom of the ethyl group reflects a significant contribution from a wagging of the entire group. The motion of the carbon and oxygen atoms which are bound to the chelate ring is considerably more restricted in a direction parallel to the bonds joining them to the ring. The motion of the oxygen atoms is
suggestive of a non-linear hydrogen bond, and this may well be the case, as both of the $\mathrm{N}-\mathrm{O}-\mathrm{O}$ angles are considerably smaller than the $104 \cdot 5^{\circ} \mathrm{H}-\mathrm{O}-\mathrm{H}$ angle in water.
Calculation of interatomic distances corrected for thermal motion assuming a riding model does not significantly change bond distances in the chelate ring, but does result in an increase of about $0.02 \AA$ in the bonds external to the ring.
As an aid to visualization of the crystal packing, Fig. 3 provides a perspective view, along $\mathbf{a}^{*}$, of a unit cell and those molecules which contribute to its contents. The molecules in adjacent layers are separated by a distance of $3.286 \AA$. The packing is such that the nickel atom of one molecule lies approximately above an oxygen atom in an adjacent molecule.
The results of this investigation and a study of nickel(II) dimethylglyoxime $\left[\mathrm{Ni}(\mathrm{DMG})_{2}\right.$ ], which will soon be reported, lead us to believe that reports of OHO hydrogen bonds with $\mathrm{O} \cdots \mathrm{O}$ distances less than about $2 \cdot 4 \AA$ should be treated with suspicion. [A $2 \cdot 42 \AA$ $\mathrm{O}--\mathrm{O}$ distance has been found in a neutron diffraction study of the related bis(2-amino-2-methyl-3-butanone oximato) nickel(II) chloride monohydrate by Schlemper, Hamilton \& La Placa (1971).] The available structural information on OHO hydrogen bonds indicates that oxygen-oxygen repulsions make improbable any shorter O---O distances.


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